Numerical simulation of incompressible flows with particle-laden interfaces

Arnold Reusken, Yuanjun Zhang*

RWTH Aachen

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2 Particle-laden fluid interfaces





Research focus

Two phase flows with transport phenomena and variable surface properties.

Methodology

- Own development; Open source (LGPL).
- Algorithms + Data structures + Post-processing.
- C++ object oriented, generic programming.
- Parallelization.

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2 Particle-laden fluid interfaces

Standard model for two-phase flows



- We assume that the interface is a sharp dividing surface.
- We assume the fluid to be incompressible, viscous, Newtonian, pure, and isothermal. Fluid dynamics: Navier-Stokes equations in separate phases.

$$\begin{cases} \rho_i \left(\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = \operatorname{div} \sigma_i + \rho_i \mathbf{g} \\ \operatorname{div} \mathbf{u} = 0 \end{cases}$$

with
$$\sigma_i = -p\mathbf{I} + \mu_i \left(\nabla \mathbf{u} + (\nabla \mathbf{u})^T \right)$$

• Standard interface condition for a clean interface:

 $[\sigma \mathbf{n}_{\Gamma}] = -\tau \kappa \mathbf{n}_{\Gamma} + \nabla_{\Gamma} \tau, \quad [\mathbf{u}]_{\Gamma} = \mathbf{0},$ $V_{\Gamma} = \mathbf{u} \cdot \mathbf{n}_{\Gamma}$

with V_{Γ} normal velocity of the interface, τ surface tension coefficient, κ curvature.

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• Interface condition represented as localized force term (constant τ).

$$f_{\mathsf{\Gamma}}(\mathsf{v}) := -\int_{\mathsf{\Gamma}} au \kappa \mathbf{n} \cdot \mathbf{v} \mathrm{d}s$$

- Level set method for interface representation.
- Finite Element methods. Extended-FEM for discontinuous quantities.
- Fast iterative solvers, multilevel techniques.
- Parallelization with MPI.

DROPS - Interface Transport Phenomena

Mass Transport

$$\frac{\partial c}{\partial t} + \mathbf{u} \cdot \nabla c = \operatorname{div}(D_i \nabla c) \quad \text{in } \Omega_i,$$
$$D_i \nabla c |_{\Gamma} \cdot \mathbf{n} = 0, \ c_1 = C_{\mu} c_2 \quad \text{on } \Gamma.$$

Henry condition : discontinuity in c. D_i : piecewise constant.

Surfactant Transport

$$\frac{\partial S}{\partial t} + \mathbf{u} \cdot \nabla S + S \operatorname{div}_{\Gamma} \mathbf{u} = \operatorname{div}_{\Gamma} (D_{\Gamma} \nabla_{\Gamma} S) \quad \text{on } \Gamma.$$

- τ does not depend on *S*.
- Observation: accumulation of surfactant at the bottom.
- For both models special numerical methods have been developed.



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Surfactant versus Solid Particles

Surfactant

Surface-active materials: bipolar large molecules with hydrophobic and hydrophilic segments.

Solid particles

Sub-micron solid particles may also display surface activity due to three-phase equilibrium contact angle.

- θ Fluid I Solid Fluid II
 - Attracted to an interface.
 - etc.



- Alter interfacial tension.
- Stabilize fluid-fluid dispersions.

Complicated microscopic mechanisms and phenomena.

- Self-assembly of particles due to long range repulsive force and attractive force.
- Formation of two-dimensional structures.
- In general, adding particles at fluid interface will alter surface properties significantly.→ Applications: Emulsion stabilization (*Pickering* emulsion), health products, etc.



E. Vignati and R. Piazza, Pickering Emulsions: Interfacial Tension, Colloidal Layer Morphology, and Trapped-Particle Motion Langmuir 2003, 19

Modeling of Particle-laden Fluid Interfaces

- We investigated continuum description of surface stress tensor of the particle-laden fluid interfaces. (Macroscopic models)
- Advantage: treatment of realistic two-phase fluids.
- Challenge: choice of effective surface coefficients (e.g. viscosity).

Boussinesq-Scriven law

$$\sigma_{\Gamma} = [\tau + (\lambda_{\Gamma} - \mu_{\Gamma}) \operatorname{div}_{\Gamma} u] \mathbf{P} + 2\mu_{\Gamma} D_{\Gamma}, \quad D_{\Gamma} = \frac{1}{2} \mathbf{P} (\nabla_{\Gamma} u + \nabla_{\Gamma} u^{T}) \mathbf{P},$$

with $\mathbf{P} = \mathbf{I} - \mathbf{nn}^{T}$ orthogonal projection, coefficients $\lambda_{\Gamma}, \mu_{\Gamma}$.

- Surface tension term.
- Dilatational viscosity term.
- Surface shear viscosity term.

Remarks

- B-S model extends the standard surface tension model ($\sigma_{\Gamma} = \tau \mathbf{P}$) by adding viscous terms.
- Effective viscosities are needed to model effects of nanoparticles.
- Not clear, yet, the range of systems this model is appropriate.

Treatment of general surface tension tensors

Interface stress tensor

• Standard surface tension model :

$$[\sigma \mathbf{n}_{\Gamma}] = - au \kappa \mathbf{n} +
abla_{\Gamma} au = \operatorname{div}_{\Gamma}(au \mathbf{P}) =: \operatorname{div}(\sigma_{\Gamma})$$

with σ_{Γ} surface stress tensor.

• General interface condition :

$$[\sigma \mathbf{n}_{\Gamma}] = \operatorname{div}_{\Gamma}(\sigma_{\Gamma})$$

• Boussinesq-Scriven law :

$$\operatorname{div}_{\Gamma}(\sigma_{\Gamma}) = \operatorname{div}_{\Gamma}(\tau \mathbf{P} + (\lambda_{\Gamma} - \mu_{\Gamma})\operatorname{div}_{\Gamma} u \mathbf{P} + 2\mu_{\Gamma}D_{\Gamma})$$

Discretization of surface force

- Improved Laplace-Beltrami discretization is applied.
- Approximation quality: $\mathcal{O}(h_{\Gamma})$

Slattery's model

$$\sigma_{\Gamma} = (1 - \mathbf{x})\tau \mathbf{P} + 2\epsilon \mathbf{D}_{\Gamma}$$

where $\epsilon = Rx(\mu_1 + \mu_2)$ is the surface shear viscosity.

Lishchuk's model

$$\mu_{\Gamma} = \frac{5}{3}(\mu_1 + \mu_2)Rx \quad \text{surface shear viscosity}$$
$$\lambda_{\Gamma} = 5(\mu_1 + \mu_2)Rx \quad \text{dilatational viscosity}$$

Note : R radius of particles, x coverage of particles on the interface, μ_i viscosity of bulk fluid.

Slattery, Sagis, Oh, Interfacial Transport Phenomena, 2nd Edition, Springer, 2007. Lishchuk, Halliday, Effective Surface Viscosities of a Particle-laden Fluid Interface, Phys. Rev. E 80, 016306(2009).

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Property	Fluid I	Fluid II
ρ	1	1
μ	1	1
au	1	
λ_{F}	1(0)	
μ_{F}	1(0)	
g	0	
v (0)	0	





- Investigation of visco-elastic interface model.
- Investigation on how to obtain effective surface coefficients.
- Improvement of iterative solver for models with viscous interface terms.
- Validation with experimental results.

More information: www.igpm.rwth-aachen.de/DROPS

• Laplace- Beltrami discretization ($\sigma_{\Gamma} = \tau \mathbf{P}$):

$$f_{\Gamma}(\mathbf{v}_{h}) = \tau \int_{\Gamma} \nabla_{\Gamma} \operatorname{id}_{\Gamma} \cdot \nabla_{\Gamma} \mathbf{v}_{h} \, ds \approx \int_{\Gamma_{h}} \nabla_{\Gamma_{h}} \operatorname{id}_{\Gamma_{h}} \cdot \nabla_{\Gamma_{h}} \mathbf{v}_{h} \, ds =: f_{\Gamma_{h}}(\mathbf{v}_{h})$$

Approximation quality : $\mathcal{O}(\sqrt{h_{\Gamma}})$ [*S.Gross*, *A.Reusken*, *SINUM*07] • This discretization can be improved:

$$\widetilde{f}_{\Gamma_h}(\mathbf{v}_h) = \tau \int_{\Gamma_h} \widetilde{\mathbf{P}}_h(\mathbf{x}) \nabla \operatorname{id}_{\Gamma_h} \cdot \nabla_{\Gamma_h} \mathbf{v}_h \, ds.$$

with $\tilde{\mathbf{n}}_h(\mathbf{x}) := \frac{\nabla \phi_h(\mathbf{x})}{\|\nabla \phi_h(\mathbf{x})\|}, \tilde{\mathbf{P}}_h(\mathbf{x}) := \mathbf{I} - \tilde{\mathbf{n}}_h(\mathbf{x})\tilde{\mathbf{n}}_h(\mathbf{x})^T$. Approximation quality : $\mathcal{O}(\mathbf{h}_{\Gamma})$ [S.Gross, A.Reusken, SINUM07]

• Similar techniques has been generalized to the case with general surface tension tensor.

