NUMERICAL SIMULATION OF LASER-INDUCED CAVITATION BUBBLES

Mathieu Bachmann

Institut für Geometrie und Praktische Mathematik, RWTH Aachen

Joint work with :

Siegfried Müller, Josef Ballmann, RWTH Aachen.

Mohsen Alizadeh, Dennis Kröninger, Thomas Kurz, Werner Lauterborn Universität Göttingen.

Philippe Helluy, Hélène Mathis, Université de Strasbourg.

Outline

- 1. Investigation in 1d
 - (a) Mathematical model
 - (b) Two different approaches : Saurel-Abgrall and the level-set function
 - (c) Initial data
 - (d) Comparison between experiments and approaches
- 2. Investigation in 2d
 - (a) Mathematical model
 - (b) Modified real ghost fluid method in 2D
 - (c) Correction due to the levelset
 - (d) Results of a spherical bubble near a rigid wall
- 3. Outlook

1d : Mathematical model

• The 1d-Euler equations in spherical coordinates

$$\frac{\partial}{\partial t}(r^2 \rho) + \frac{\partial}{\partial r}(r^2(\rho v_r)) = 0$$

$$\frac{\partial}{\partial t}(r^2 \rho v_r) + \frac{\partial}{\partial r}(r^2(\rho v_r^2 + p)) = 2 p r$$

$$\frac{\partial}{\partial t}(r^2 \rho E) + \frac{\partial}{\partial r}(r^2(\rho v_r(E + p/\rho))) = 0$$

• The stiffened gas pressure law is used to close the system. $p(\rho, e, \varphi) = (\gamma(\varphi) - 1)\rho e - \gamma(\varphi)\pi(\varphi).$ (1)

 φ is the phase indicator function (gas fraction, level set function).

Saurel Abgrall Approach

• The two phases (gas and liquid) are distinguished by the mass fraction φ which satisfies a transport equation without mass transfer.

$$\frac{\partial \varphi}{\partial t} + v_r \frac{\partial \varphi}{\partial r} = 0.$$

- \bullet For the pure phases, the coefficients γ and π are obtained by measurements.
- A linear interpolation between the two phases is used for the mixture,

$$\beta_1(\varphi) = \varphi \beta_1(1) + (1 - \varphi) \beta_1(0),$$

$$\beta_2(\varphi) = \varphi \beta_2(1) + (1 - \varphi) \beta_2(0).$$

where β_1 and β_2 are defined by $\beta_1 = 1/(\gamma - 1)$ and $\beta_2 = \gamma \pi/(\gamma - 1)$.

Real Ghost Fluid Method (Wang, Liu, Khoo)

• The two phases are distinguished by the function levelset (distance function) which satisfies a transport equation.

$$\partial_t \varphi + \mathbf{v} \cdot \nabla \varphi = 0.$$

Fluid A Ghost Fluid A i-2 i-1 i i+1 i+2 • A Riemann problem is defined at the interface \mathbf{P}_{1} $U_{\mathbf{R}} = U_{l+2}$ $(\varphi = 0)$ and solved for 11. $\mathbf{U}_{\mathrm{L}} = \mathbf{U}_{\mathrm{L}}$ PIT PIR predicting the interfacial states (ρ_{IL} , ρ_{IR} , p_I and i-2 1-1 **i+1** i+2 u_I). Ghost Fluid B Fhuid B Phase boundary

Initial Data : Fitting of Equilibrium Radius

Initial conditions :

- $t_{max} = 70.7 \ \mu s$ ("Exp")
- $\mathsf{R}_b = \mathsf{R}_{max}$ (Exp)
- $\dot{R}_b = 0$

 $\Rightarrow R_{eq} = 6.92 \times 10^{-5}$ m in minimizing the least square error of the Keller-Miksis model.



	Initial data		Material parameters				
	$ ho~[{ m kg/m}^3]$	p [Pa]	γ [-]	π [Pa]	$c_v \; [{J}/{kg} \; {K}]$	${\cal R}~[{\sf J}/{\sf kg}~{\sf K}]$	
Gas	9.5e-4	4.57	1.4	0	708.3	283.32	
Liquid	1000	100000	1.1	2.e+9	4190.0	418	

7

Numerical Results: $R_{eq} = 6.92 \times 10^{-5} m$



Numerical Results: $R_{eq} = 6.92 \times 10^{-4} m$

	Initial data		Material parameters				
	$ ho~[{ m kg/m}^3]$	$p \; [Pa]$	γ [-]	π [Pa]	$c_v \; [{J}/{kg} \; {K}]$	${\cal R}~[{\sf J}/{\sf kg}~{\sf K}]$	
Gas	9.57e-1	72560	1.4	0	708.3	283.32	
Liquid	1000	100000	1.1	2.e+9	4190.0	418	



8

2d : Mathematical model

• The 2d-Euler equations with rotational symmetry

$$\begin{aligned} \frac{\partial}{\partial t}(\rho) &+ \frac{\partial}{\partial r}(\rho v_r) + \frac{\partial}{\partial z}(\rho v_z) + \frac{1}{r}\rho v_r = 0\\ \frac{\partial}{\partial t}(\rho v_r) &+ \frac{\partial}{\partial r}(\rho v_r^2 + p) + \frac{\partial}{\partial z}(\rho v_r v_z) + \frac{1}{r}\rho v_r^2 = 0\\ \frac{\partial}{\partial t}(\rho v_z) &+ \frac{\partial}{\partial r}(\rho v_r v_z) + \frac{\partial}{\partial z}(\rho v_z^2 + p) + \frac{1}{r}\rho v_z v_r = 0\\ \frac{\partial}{\partial t}(\rho E) &+ \frac{\partial}{\partial r}(v_r(\rho E + p)) + \frac{\partial}{\partial z}(v_z(\rho E + p)) + \frac{1}{r}v_r(\rho E + p) = 0\end{aligned}$$

• The stiffened gas pressure law is used to close the system. $p(\rho, e, \varphi) = (\gamma(\varphi) - 1)\rho e - \gamma(\varphi)\pi(\varphi)$ with φ the level-set function.

Modified Real Ghost Fluid Method

- The levelset defines the normal to the interface $N = \frac{\nabla \varphi}{|\nabla \varphi|}$
 - A Riemann problem is defined at the interface between the cell A and B such that the angle made by the respective normals are the minimum
 - The interfacial states define the ghost cells (ρ_{IL} , ρ_{IR} , p_I and u_I).



- A single-phase Riemann problem is solved at cell interfaces with ghost cells as boundary condition.
- The solution can be advanced to the next time step.

Change of the levelset sign in a cell after update

• Modification of the energy using the corresponding gas law.

$$\mathbf{U}_{\mathbf{k}}^{n+1} = (\varrho, \varrho v, \varrho E) \to p(\varrho, e, EOS(\phi_{k}^{n})) \to p(\varrho, p, EOS(\phi_{k}^{n+1}))$$
$$\to \tilde{\mathbf{U}}_{\mathbf{k}}^{n+1} = \left(\varrho, \varrho v, \varrho \tilde{E}\right).$$

• Modification of the density with the corresponding ghost cell.

$$\tilde{\mathbf{U}}_{\mathbf{k}}^{n+1} = \left(\varrho, \varrho v, \varrho \tilde{E}\right) \to \tilde{\mathbf{U}}_{\mathbf{k}}^{n+1} = \left(\varrho^{GC}, \varrho^{GC} v, \varrho^{GC} \tilde{E}^{GC}\right).$$

Results : Spherical bubble near a rigid wall



• Initial data

	Initial data		Material parameters			
	$ ho~[{ m kg/m}^3]$	p [Pa]	γ [-]	π [Pa]	$c_v \; [{\sf J}/{\sf kg} \; {\sf K}]$	${\cal R}~[{\sf J}/{\sf kg}~{\sf K}]$
Gas	0.026077	2118	1.4	0	717.5	283.32
Liquid	1000	50000000	7.15	3.e+8	201.1	1236.765

12

Results : Spherical bubble near a rigid wall

6 7 8 9 10 11 12 13 14 15 16 17 1







 $t=1.59~\mu s$

 $t = 3.25 \ \mu s$

0.002 x [m] 0.003

1

0.004

Aachen, February 2009



Aachen, February 2009

 $\phi = 0$

0.0018

0.0016

log(1+|∇p|): 1 2 3 4 5 6 7 8 9 10 11 12 13 14 15 16 17 18





x [m]

0.0014

0.0012

t= 8.20 μ s

15

Aachen, February 2009



Aachen, February 2009



Aachen, February 2009



Results : Spherical bubble near a rigid wall



Extract along the symmetry axis



Conclusion

Saurel-Abgrall:

Real Ghost Fluid Method:

- Severe numerical phase transition No phase transition
- Rebound overpredicted
 Re
 - Rebound well-predicted

Real Ghost Fluid Method in 2D:

- Difficulties to maintain the interface sharp due to the geometry and the extreme conditions
- Loss of the gas phase

Future Work

- Investigation of shock bubble interactions.
- Qualitative comparisons with experimental data.